

Big Data Analytics Challenges

Data Analytics = Information Compression

- □ From 23 ... to ... 2^{23} (10M) $\left(\underbrace{23}_{2 \#' s} \to \underbrace{2^{23}}_{8 \#' s}\right)$
- ☐ Two centuries of Data Science: 1798 → 2020
- □ In the 18th century, Henry Cavendish used just 23 observations to answer a fundamental question "What is the Mass of the Earth?" He estimated very accurately the mean density of the Earth/H₂O (5.483±0.1904 g/cm³)
- □ In the 21st century to achieve the same scientific impact, matching the reliability and the precision of the Cavendish's 18th century prediction, requires a monumental community effort using massive and complex information often exceeding 10M (2²³) bytes

Dinov (2016) J MedicalStat



Common Characteristics of Big Data

IBM Big Data 4V's: Volume, Variety, Velocity & Veracity

Big Bio Data Dimensions	Tools
Size	Harvesting and management of vast amounts of data
Complexity	Wranglers for dealing with heterogeneous data
Incongruency	Tools for data harmonization and aggregation
Multi-source	Transfer and joint multivariate representation & modeling
Multi-scale	Macro → meso → micro → nano scale observations
Time	Techniques accounting for longitudinal effects (e.g., time corr)
Incomplete	Reliable management of missing data, imputation

Example: analyzing observational data of 1,000's Parkinson's disease patients based on 10,000's signature biomarkers derived from multi-source imaging, genetics, clinical, physiologic, phenomics and demographic data elements

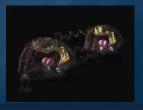
Software developments, student training, service platforms and methodological advances associated with the Big Data Discovery Science all present existing opportunities for learners, educators, researchers, practitioners and policy makers

Dinov, GigaScience (2016) PMID:26918190

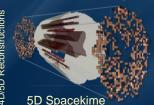


Longitudinal Data Analytics

- <u>Neuroimaging</u>:
 - □ 4D fMRI: time-series, represents measurements of hydrogen atom densities over a 3D lattice of spatial locations ($1 \le x, y, z \le 64$ pixels), about 3×3 millimeters² resolution. Data is recorded longitudinally over time ($1 \le t \le 180$) in increments of about 3 seconds, then post-processed
 - ☐ State-of-the-art Approaches: Time-series modeling or Network analysis
 - □ Spacekime Analytics: 5D fMRI kime-series, represent the hydrogen atom densities over the same 3D lattice of spatial locations, longitudinally over the 2D kime space, $\kappa = re^{i\phi} \in \mathbb{C}$
 - ☐ *Differences*: Spacekime analytics estimate and utilize the kime-phases

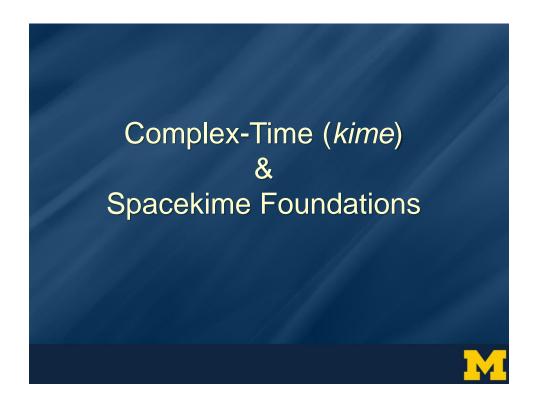


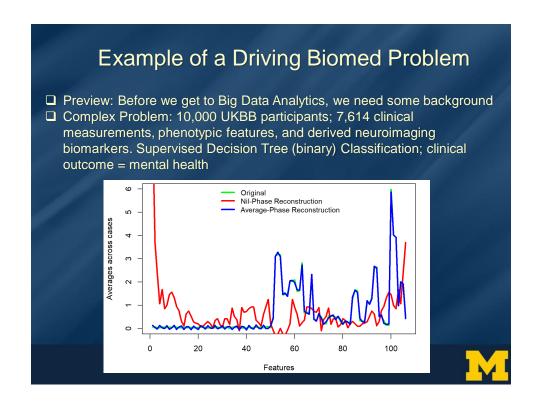




4D Spacetime







The Fourier Transform

By separability, the classical **spacetime Fourier transform** is just four Fourier transforms, one for each of the four spacetime dimensions, (x,t)=(x,y,z,t). The FT is a function of the <u>angular frequency</u> ω that propagates in the wave number direction k (<u>space frequency</u>). Symbolically, the forward and inverse Fourier transforms of a 4D (n=4) spacetime function f, are defined by:

$$FT(f) = \hat{f}(\mathbf{k}, \omega) = \frac{1}{(2\pi)^{\frac{n}{2}}} \int f(\mathbf{x}, t) e^{i(\omega t - \mathbf{k}\mathbf{x})} dt d^3 \mathbf{x},$$

$$IFT(\hat{f}) = \hat{f}(\mathbf{x}, t) = \frac{1}{(2\pi)^{\frac{n}{2}}} \int \hat{f}(\mathbf{k}, \omega) e^{-i(\omega t - \mathbf{k}\mathbf{x})} d\omega d^3 \mathbf{k}.$$

$$\left[\hat{f}(\mathbf{x},t) = IFT(\hat{f}) = IFT(FT(f)) = f(\mathbf{x},t), \quad \forall \mathbf{z} \in \mathbb{C}, \mathbf{z} = \underbrace{A}_{mag} e^{i \underbrace{\varphi}_{phase}}\right]$$



Fourier Came Festel About Nelp Sazeshot Wumber of Terms 1 100 Figure Repeted 1 120 16 Magnitudes Final Receity Four Clear Final Receity Four Receiver F

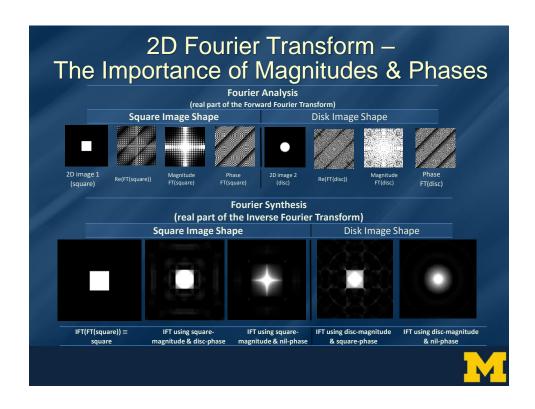
1D Fourier Transform Example

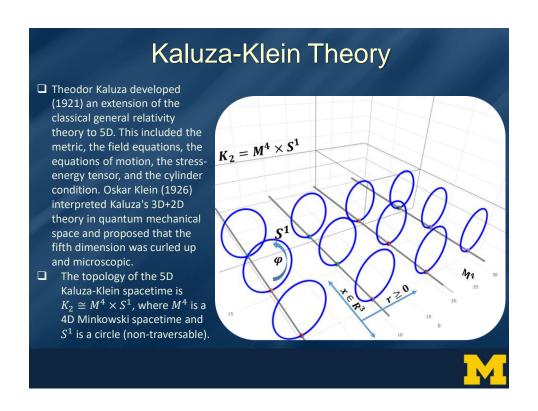
SOCR 1D Fourier / Wavelet signal decomposition into magnitudes and phases (Java applet)

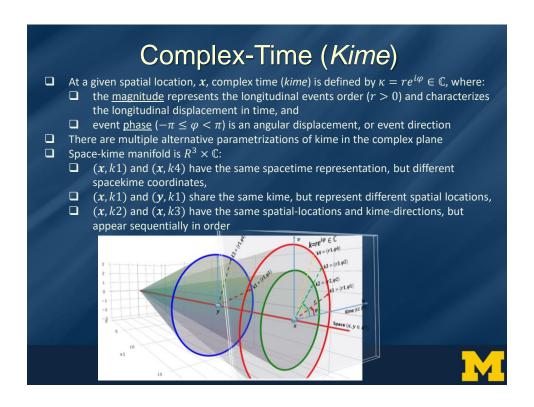
<u>Top-panel</u>: original signal (image), white-color curve drawn manually by the user and the reconstructed synthesized (IFT) signal, red-color curve, computed using the user modified magnitudes and phases <u>Bottom-panels</u>: the Fourier analyzed signal (FT) with its magnitudes and phases

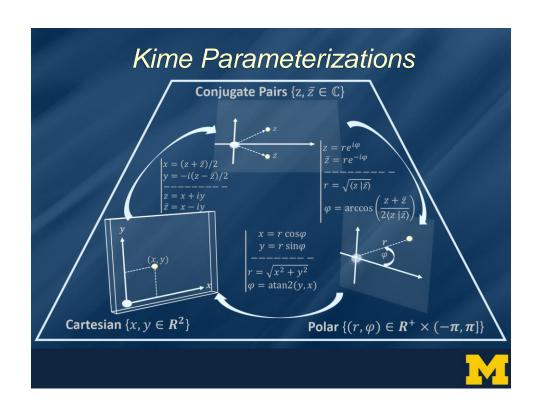
http://www.socr.ucla.edu/htmls/game/Fourier_Game.html (Java Applet)











The Spacekime Manifold

- Spacekime: $(x, k) = \left(\underbrace{x^1, x^2, x^3}_{\text{space}}, \underbrace{c\kappa_1 = x^4, c\kappa_2 = x^5}_{\text{kime}}\right) \in X, \ c \sim 3 \times 10^8 \ m/s$
- **Kevents** (complex events): points (or states) in the spacekime manifold X. Each kevent is defined by where (x = (x, y, z)) it occurs in space, what is its causal longitudinal order $(r = \sqrt{(x^4)^2 + (x^5)^2})$, and in what kime-direction $(\varphi = a \tan 2(x^5, x^4))$ it takes place.
- Spacekime interval (ds) is defined using the general Minkowski 5×5 metric tensor $(\lambda_{ij})_{i=1,j=1}^{5,5}$, which characterizes the geometry of the (generally curved) spacekime manifold:

 $ds^2 = \sum_{i=1}^{5} \sum_{j=1}^{5} \lambda_{ij} dx^i dx^j = \lambda_{ij} dx^i dx^j$ $(\lambda_{ij}) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix}$ Euclidean (flat) spacekime metric corresponds to the tensor:

- \square Spacelike intervals correspond to $ds^2 > 0$, where an inertial frame can be found such that two kevents $a, b \in X$ are simultaneous. An object can't be present at two kevents which are separated by a spacelike interval.
- \Box Lightlike intervals correspond to $ds^2 = 0$. If two kevents are on the line of a photon, then they are separated by a lightlike interval and a ray of light could travel between the two kevents.
- <u>Kimelike</u> intervals correspond to $ds^2 < 0$. An object can be present at two different kevents, which are separated by a kimelike interval.



Spacekime Calculus

□ Kime Wirtinger derivative (first order kime-derivative at
$$k = (r, φ)$$
):
$$f'(z) = \frac{\partial f(z)}{\partial z} = \frac{1}{2} \left(\frac{\partial f}{\partial x} - i \frac{\partial f}{\partial y} \right) \text{ and } f'(\bar{z}) = \frac{\partial f(\bar{z})}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial f}{\partial x} + i \frac{\partial f}{\partial y} \right).$$

 $df = \partial f + \bar{\partial} f = \frac{\partial f}{\partial z} dz + \frac{\partial}{\partial \bar{z}} d\bar{z}$ In Conjugate-pair basis:

In Polar kime coordinates:

$$f'(k) = \frac{\partial f(k)}{\partial k} = \frac{1}{2} \left(\cos \varphi \frac{\partial f}{\partial r} - r \sin \varphi \frac{\partial f}{\partial \varphi} - \mathbf{I} \left(\sin \varphi \frac{\partial f}{\partial r} + \frac{1}{r} \cos \varphi \frac{\partial f}{\partial \varphi} \right) \right)$$

$$f'(\bar{k}) = \frac{\partial f(\bar{k})}{\partial \bar{k}} = \frac{1}{2} \left(\cos \varphi \, \frac{\partial f}{\partial r} - r \, \sin \varphi \, \frac{\partial f}{\partial \varphi} + \frac{1}{\ell} \left(\sin \varphi \, \frac{\partial f}{\partial r} + \frac{1}{r} \cos \varphi \, \frac{\partial f}{\partial \varphi} \right) \right).$$

 \square Kime Wirtinger acceleration (second order kime-derivative at $k = (r, \varphi)$):

$$f''(\mathbf{k}) = \frac{1}{4r^2} \left((\cos \varphi - \mathbf{1} \sin \varphi)^2 \left(2\mathbf{1} \frac{\partial f}{\partial \varphi} - \frac{\partial^2 f}{\partial \varphi^2} + r \left(-\frac{\partial f}{\partial r} - 2\mathbf{1} \frac{\partial^2 f}{\partial r \partial \varphi} + r \frac{\partial^2 f}{\partial r^2} \right) \right) \right).$$



Spacekime Calculus

☐ Kime Wirtinger integration:

The path-integral of a complex function $f:\mathbb{C}\to\mathbb{C}$ on a specific path connecting $z_a\in\mathbb{C}$ to $z_b \in \mathbb{C}$ is defined by generalizing Riemann sums:

$$\lim_{|z_{i+1}-z_i|\to 0} \sum_{i=1}^{n-1} (f(z_i)(z_{i+1}-z_i)) \cong \oint_{z_a}^{z_b} f(z_i)dz.$$

This assumes the path is a polygonal arc joining z_a to z_b , via $z_1=\overline{z_a},z_2,z_3,...,z_n=\overline{z_b},$ and we integrate the piecewise constant function $f(z_i)$ on the arc joining $z_i \to z_{i+1}$. Assumptions: the path $z_a \rightarrow z_b$ needs to be defined and the limit of the generalized Riemann sums, as $n \to \infty$, will yield a complex number representing the Wirtinger integral of the function over the path.

☐ Similarly, extend the classical area integrals, indefinite integral, and Laplacian:

Definite area integral: for $\Omega \subseteq \mathbb{C}$, $\int_{\Omega} f(z)dzd\bar{z}$.

Indefinite integral: $\int f(z)dzd\bar{z}, \ df = \frac{\partial f}{\partial z}dz + \frac{\partial f}{\partial \bar{z}}d\bar{z}$.

The *Laplacian* in terms of conjugate pair coordinates is $\Delta f = d^2 f = 4 \frac{\partial f}{\partial z} \frac{\partial f}{\partial \bar{z}} = 4 \frac{\partial f}{\partial \bar{z}} \frac{\partial f}{\partial z}$.

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Newton's equations of motion in kime

$$\begin{vmatrix} v = at + v_o \\ x = x_o + v_o t + \frac{1}{2}at^2 \Rightarrow \\ v^2 = 2ax + v_o^2 \end{vmatrix} \Rightarrow \begin{vmatrix} v = a_1k_1 + v_{o1} = a_2k_2 + v_{o2} \\ x = x_{o1} + v_{o1}k_1 + \frac{1}{2}a_1k_1^2 = x_{o2} + v_{o2}k_2 + \frac{1}{2}a_2k_2^2 \\ v v_1 = 2a_1x + v_{o1}^2 \\ v v_2 = 2a_2x + v_{o2}^2 \end{vmatrix}$$

- Derived from the Kime Wirtinger velocity and acceleration
 - \square Kime-velocity ($\mathbf{k}=(t,\varphi)$) is defined by the Wirtinger derivative of the position with respect to kime:

$$\nu(\mathbf{k}) = \frac{\partial \mathbf{x}}{\partial \mathbf{k}} = \frac{1}{2} \left(\cos \varphi \frac{\partial \mathbf{x}}{\partial t} - \frac{1}{t} \sin \varphi \frac{\partial \mathbf{x}}{\partial \varphi} - i \left(\sin \varphi \frac{\partial \mathbf{x}}{\partial t} + \frac{1}{t} \cos \varphi \frac{\partial \mathbf{x}}{\partial \varphi} \right) \right)$$

$$v_2 = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{\sqrt{dx^2 + dy^2}} = \frac{\sqrt{dx^2 + dy^2}}{\sqrt{dx^2 + dy^2}}$$

- □ The directional kime derivatives v_1 and v_2 : $v_1 = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{dk_1} = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{\cos(\varphi) \, dt t \sin(\varphi) \, d\varphi}, \qquad v_2 = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{dk_2} = \frac{\sqrt{dx^2 + dy^2 + dz^2}}{\sin(\varphi) \, dt + t \cos(\varphi) \, d\varphi}.$ □ Equ 1: As $a_1 = \frac{\partial v}{\partial k_1}$ and $a_2 = \frac{\partial v}{\partial k_2}$, integrating both sides yields $\int a_1 \, dk_1 = \int dv \, and \int a_2 \, dk_2 = \int dv$. Since the acceleration is constant in kime, $v = \frac{\partial v}{\partial k_2}$ and $v = \frac{\partial v}{\partial k_2}$ and $v = \frac{\partial v}{\partial k_2}$. $a_1 \int dk_1 = a_1 k_1 + v_{o1} = a_2 \int dk_2 = a_2 k_2 + v_{o2}$, where v_{o1} and v_{o2} are constants representing the initial k-velocities, defined in relation to the
- kime dimensions k_1 and k_2 , respectively. $\Box \quad \underline{\operatorname{Equ} 2} : v = \frac{\partial f_1}{\partial k_1} = a_1 k_1 + v_{o1} \text{ and } v = \frac{\partial f_2}{\partial k_2} = a_2 k_2 + v_{o2} \text{ (from equ 1), integrating we get } \int \partial f_1 = \int a_1 k_1 \partial k_1 + v_{o1} \int \partial k_1 \text{ and } \int \partial f_2 = \int a_2 k_2 \partial k_2 + v_{o2} \int \partial k_1 \partial k_1 \partial k_2 + v_{o3} \partial k_2 \partial k_2 + v_{o4} \partial k_2 \partial k_2 \partial k_3 \partial k_4 \partial k_4 \partial k_4 \partial k_4 \partial k_4 \partial k_4 \partial k_5 \partial k_4 \partial k_5 \partial k_4 \partial k_5 \partial k_5$ $v_{o2}\int \partial k_2$. As a_1 and a_2 are constants, we have $x=a_1\int \partial \frac{k_1^2}{2}+v_{o1}k_1=a_1\frac{k_1^2}{2}+v_{o1}k_1+C_1$ and we can compute the constant $C_1=x_{o1}$ by setting $k_1=0$. Analogously, we will have $x=a_2\int \partial \frac{k_2^2}{2}+v_{o2}k_2=a_2\frac{k_2^2}{2}+v_{o2}k_2+C_2$, and we estimate the constant $C_2=x_{o2}$ by setting $k_2=0$.
- $\Box \quad \underline{\text{Equ 3:}} \ a_1 = \frac{\partial v}{\partial k_1} = \frac{\partial v}{\partial x} \times \frac{\partial x}{\partial k_1} = \frac{\partial v}{\partial x} \times v_1 = v_1 \frac{\partial v}{\partial x} \text{ Again integrating, we get } \int a_1 dx = \int v_1 dv = \int \frac{v}{\cos(\phi)} dv = \frac{1}{\cos(\phi)} \int v dv \text{ and thus, } a_1 x + C_1 = \frac{\partial v}{\partial x} + \frac{\partial v}{\partial x}$ $\frac{v^2}{2\cos(\varphi)}. \text{ Under the initial condition } (v_o=v(0)) \text{ this becomes } 2a_1x+v_{o1}^2=vv_1.$ $\boxed{\quad \underline{\text{Equ 4:}}} \text{ Analogously, we will have } 2a_2x+v_{o2}^2=vv_2.$



Spacekime Generalizations

 \square Spacekime generalization of <u>Lorentz transform</u> between two reference frames, K & K':

(the interval ds is Lorentz transform invariant)

$$\underbrace{\begin{pmatrix} x' \\ y' \\ z' \\ k'_1 \\ k'_2 \end{pmatrix}}_{\in K'} = \begin{pmatrix} \zeta & 0 & 0 & -\frac{c^2}{v_1}\beta^2\zeta & -\frac{c^2}{v_2}\beta^2\zeta \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ -\frac{1}{v_1}\beta^2\zeta & 0 & 0 & 1 + (\zeta - 1)\frac{c^2}{(v_1)^2}\beta^2 & (\zeta - 1)\frac{c^2}{v_1v_2}\beta^2 \\ -\frac{1}{v_2}\beta^2\zeta & 0 & 0 & (\zeta - 1)\frac{c^2}{v_1v_2}\beta^2 & 1 + (\zeta - 1)\frac{c^2}{(v_2)^2}\beta^2 \end{pmatrix} \underbrace{\begin{pmatrix} x \\ y \\ z \\ k_1 \\ k_2 \end{pmatrix}}_{\in K}$$

$$\text{where}\quad 0\leq\beta=\frac{1}{\sqrt{\left(\frac{c}{\nu_1}\right)^2+\left(\frac{c}{\nu_2}\right)^2}}\leq 1\quad \&\quad \zeta=\frac{1}{\sqrt{1-\beta^2}}\geq 1\;.$$

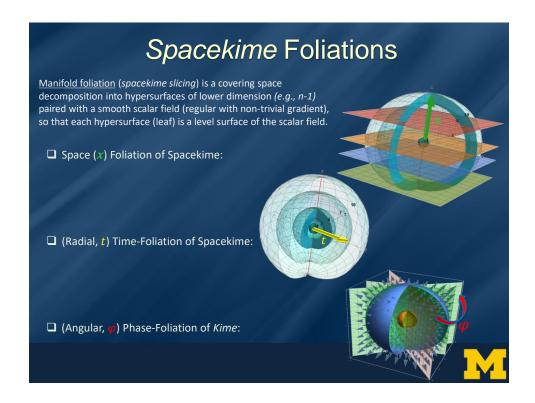
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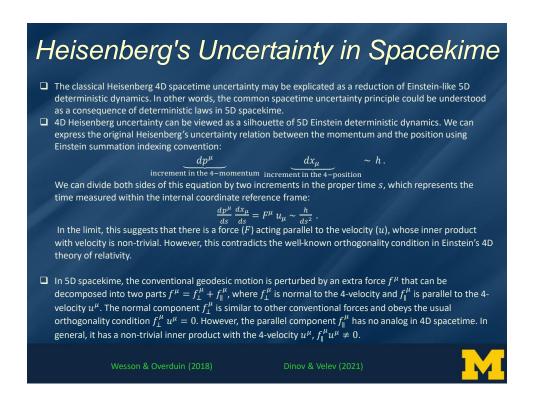


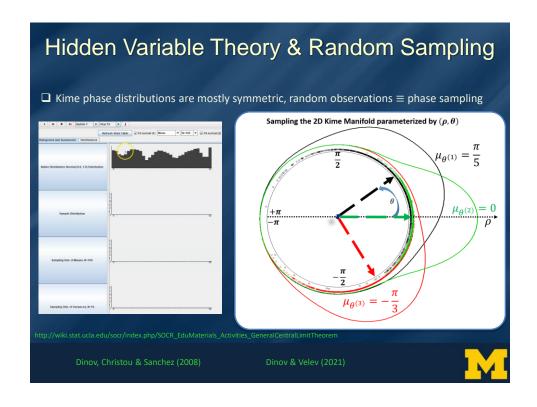
Spacekime Math Generalizations

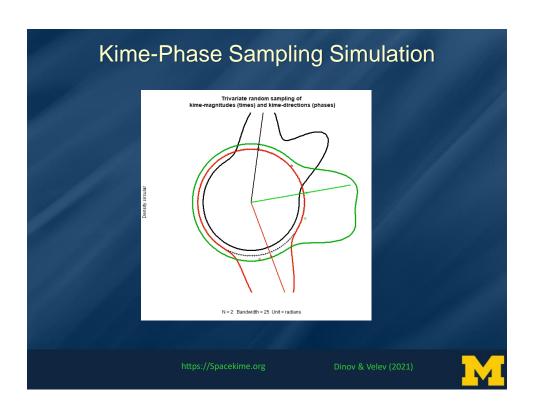
☐ Derived other spacekime concepts: law of addition of velocities, energy-momentum conservation law, stability conditions for particles moving in spacekime, conditions for nonzero rest particle mass, and causal structure of spacekime ...













 \Box Copenhagen Interpretation An instant measurement causes the wavefunction Ψ to randomly collapse only into one of the eigenfunctions of the quantity that is being measured.

$$\underbrace{ \psi = \sum_{\alpha} c_{\alpha} \psi_{\alpha} \ \, (\text{wavefunction ollapse} }_{\text{Natural, intrinsic, fuzzy}} \underbrace{ \begin{array}{c} \text{Copenhagen Interpretation} \\ \text{Measurement process} \\ \text{observe the total energy} \end{array} }_{\text{Observe the total energy}} \underbrace{ \begin{array}{c} \text{Wavefunction collapse} \\ \text{$\psi = c_{\alpha_0} \psi_{\alpha_0}$, for some index α_0,} \\ c_{\alpha_0} \in \mathbb{C} \\ \text{Total Energy} = E_{\alpha_0} \\ \text{Instance of the observed state of the system} \\ \end{array} }$$

☐ Spacekime Interpretations

$$\frac{\Psi(t)}{\text{Spacetime}} = \int_{-\pi}^{\pi} \frac{\Psi(t, \phi)}{\text{Spacekime}} \, d\phi$$

$$\text{Vunknown Total Energy} \\ \text{Indural, intrinsic, fuzzy, probabilistic} \\ \text{state of the system}$$

$$\frac{\Phi(t)}{\text{Spacekime}} = \frac{\Phi(t, \phi)}{\text{Measurement process}} \, d\phi$$

$$\text{Spacekime Interpretation} \\ \text{Measurement process} \\ \text{observe energy} \\ \text{at time } t_o$$

$$\text{Observed Total Energy} = E_{\phi'_o}, \quad \text{which} \\ \text{still represents an eigenvalue of the } \hat{H}$$

For a fixed-time instantaneous measurement of the system at $t = t_{n}$, the wavefunction, or inference-function, $\Psi(x) = \Psi(t_{n}, \varphi)$ is naturally an aggregate measure over the entire kime-phase distribution with a range $-\pi \le \varphi < \pi$. So the entire kime phases distribution (φ) may not be einerth, positionally, and instantaneously observed, the actual measurement, or inference, only reflects a measurement $\Psi(t_{n}, \varphi)$ for one random phase, φ_{n} in other words, the natural state of the system is theoretically described by a compositional configuration.

Notework special wavefunction reflects the value at a given time point $(t_i, p(s_i) = \Psi(t_i, \varphi_i))$. Where the actual observation monifests as an immutable instantaneous measurement value, $\Psi(t_i) = \Psi(t_i, \varphi_i)$. In a measure-theoretic sense, a pair of simultaneous $(t = t_i)$ independent measurements of the exact same spacekime system would naturally yield two distinct observed values: $\Psi' \equiv \Psi(t_i, \varphi_i)$ and $\Psi'' \equiv \Psi(t_i, \varphi_i)$ where the two phases are independently sampled from the curalor phase distinction, i.e., $\psi_i \in \Psi' = \Psi(t_i, \varphi_i)$. Where the two

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Spacekime Open Math Problems

- ☐ Does kime have the same interpretation in quantum mechanics and in general relativity (relative to a specified origin), just like the spatial references? In other words, is kime universal and absolute?
- ☐ We know time, by itself, is excluded from the Wheeler-DeWitt equation. Is this true for kime as well? That is, does the Wheeler-DeWitt equation depend on kime the same way it depends on the particle location?
- ☐ Is there kime-dilation, reminiscent of time-dilation? In other words, does the action of moving objects affect (slow) kime? How?
- □ Explore the relations between various spacekime principles (e.g., space-kime motion and PDEs with respect to kime) and Painlevé equations in the complex plane.
- ☐ Extend the concepts of time-based evolution, time-varying processes, and probability to the 2D kime manifold.



Spacekime Open Math Problems

☐ **Ergodicity**

Let's look at the particle velocities in the 4D Minkowski spacetime (X), a measure space where gas particles move spatially and evolve longitudinally in time. Let $\mu=\mu_X$ be a measure on X, $f(x,t)\in L^1(X,\mu)$ be an integrable function (e.g., velocity of a particle), and $T\colon X\to X$ be a measure-preserving transformation at position $x\in\mathbb{R}^3$ at time $t\in\mathbb{R}^+$.

Prove a pointwise ergodic theorem arguing that in a measure theoretic sense, the average of f over all particles in the gas system at a fixed time, $\bar{f}=E_t(f)=\int_{\mathbb{R}^3}f(x,t)d\mu_x$, will be equal to the average velocity of just one particle over the entire time span,

$$\hat{f} = \lim_{n \to \infty} \left(\frac{1}{n} \sum_{i=0}^{n} f(T^{i} x) \right)$$
. That is, prove that $\bar{f} \equiv \hat{f}$.

The spatial probability measure is denoted by μ_x and the transformation T^ix represents the dynamics (time evolution) of the particle starting with an initial spatial location $T^ox = x$. Investigate the ergodic properties of various transformations in the 5D Minkowski spacekime.

$$\bar{f} = E_t(f) = \int_{-\pi}^{+\pi} f(t,\phi) d\Phi \stackrel{?}{\cong} \lim_{t \to \infty} \left(\frac{1}{t} \sum_{i=0}^t f(t,\phi_o) \right) = \hat{f}$$

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Spacekime Open Math Problems

☐ Inference Inner Product

Define the inner product between two inference functions, $\langle \psi | \phi \rangle \equiv \langle \psi, \phi \rangle$, as a measure of the level of inference overlap, result consistency, agreement or synergies between their corresponding inferential states. The inner product provides the foundation for a probabilistic interpretation of data science inference in terms of transition probabilities. The squared modulus of an inference function, $\langle \psi | \psi \rangle = \|\psi\|^2$, represents the probability density that allows us to measure specific inferential outcomes for a given set of observables. To facilitate probability interpretation, the law of total probability requires the normalization condition, i.e., $1 = \int \|\psi\|^2$. Let's illustrate the modulus in the scope of logistic inference; the square modulus of the inference function is:

$$\begin{aligned} \|\psi\|^2 &= \langle \psi|\,\psi \rangle = \langle \psi(X,Y)|\psi(X,Y) \rangle = \langle \hat{\beta}^{oLS}|\,\hat{\beta}^{oLS} \rangle = \\ &= \langle (X^TX)^{-1}X^TY|(X^TX)^{-1}X^TY \rangle = \left((X^TX)^{-1}X^TY \right)^T (X^TX)^{-1}X^TY = \\ &= Y^TX(X^TX)^{-1}(X^TX)^{-1}X^TY = Y^T\underbrace{X(X^TX)^{-2}X^T}_{\hat{D}}Y = Y^TDY = \left\langle \left(D\frac{1}{2}\right)^TY \middle| \left(D\frac{1}{2}\right)Y \right\rangle = \|Y\|_D^2. \end{aligned}$$

What would be the effect of exploring the use of the matrix D as a constant normalization factor $(D^{\frac{1}{2}})$? Define an appropriate <u>coherence metric</u> that captures the agreement, or overlap, between a pair of complementary inference functions or data analytic strategies. E.g., inference consistency measures may be based on:

Coherence =
$$\frac{\langle \psi | \phi \rangle}{\sqrt{\langle \psi | \psi \rangle} \times \langle \phi | \phi \rangle} = \frac{\langle \psi | \phi \rangle}{\|\psi \| \|\phi \|}$$

Alternatively, as the data represent random variables (vectors, or tensors) and the specific data-analytic strategy yields the inference function, explore <u>mutual information of operators</u>, i.e., linear or non-linear operator acting on the data:

$$I(\psi; \phi) = \sum_{i} \sum_{j} \langle \psi_{i} | \phi_{j} \rangle \log \left(\frac{\langle \psi_{i} | \phi_{j} \rangle}{\|\psi_{i}\| \|\phi_{j}\|} \right),$$

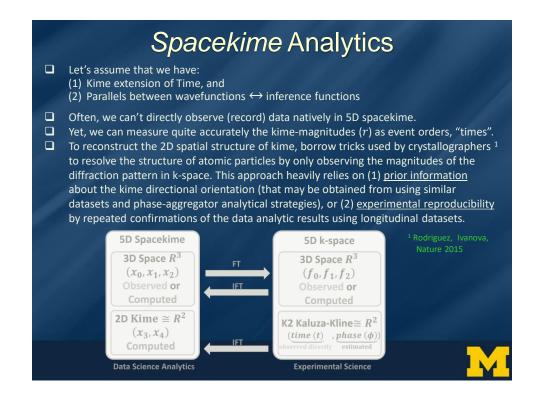
where the inference states ψ_i and ϕ_i are eigenfunctions corresponding to some eigenvalues θ_i .

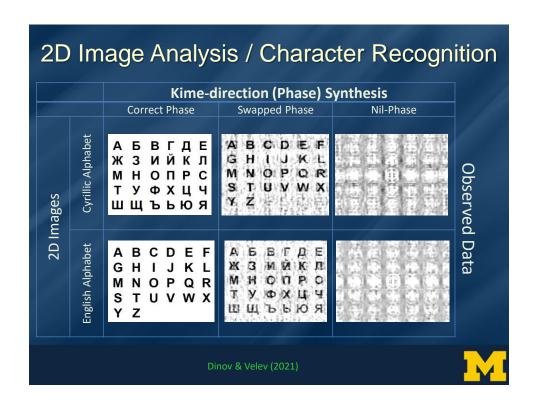


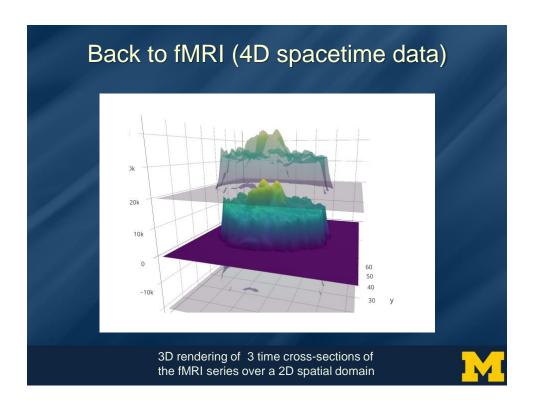


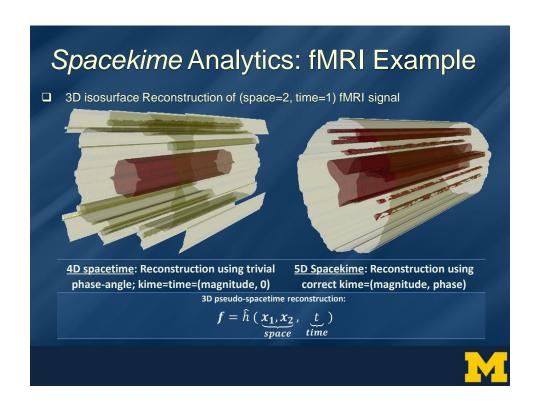
Mathematical-Physics	Data Science		
A particle is a small localized object that	An object is something that exists by itself, actually or		
permits observations and characterization of	potentially, concretely or abstractly, physically or		
its physical or chemical properties	incorporeal (e.g., person, subject, etc.)		
An <u>observable</u> a dynamic variable about	A feature is a dynamic variable or an attribute about ar		
particles that can be measured	object that can be measured		
Particle <u>state</u> is an observable particle	Datum is an observed quantitative or qualitative value,		
characteristic (e.g., position, momentum)	an instantiation, of a feature		
Particle <u>system</u> is a collection of	Problem, aka Data System, is a collection of		
independent particles and observable	independent objects and features, without necessarily		
characteristics, in a closed system	being associated with apriori hypotheses		
Wave-function	Inference-function		
Reference-Frame <u>transforms</u> (e.g., Lorentz)	Data <u>transformations</u> (e.g., wrangling, log-transform)		
State of a system is an observed	Dataset (data) is an observed instance of a set of		
measurement of all particles ~ wavefunction	datum elements about the problem system, $0 = \{X, Y\}$		
A particle system is computable if (1) the	Computable data object is a very special		
entire system is logical, consistent, complete	representation of a dataset which allows direct		
and (2) the unknown internal states of the	application of computational processing, modeling,		
system don't influence the computation	analytics, or inference based on the observed dataset		
(wavefunction, intervals, probabilities, etc.)	//		

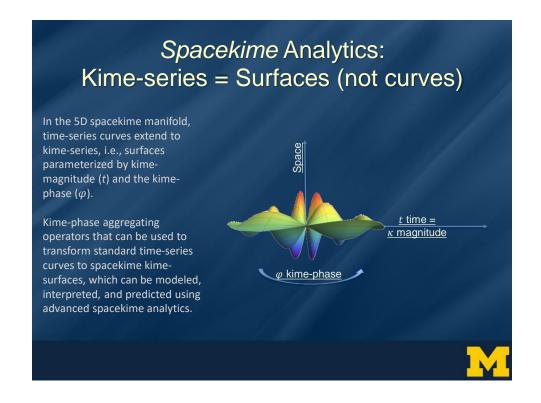
Mathematical-Physics ⇒ Data Science					
Math-Physics	Data Science				
Wavefunction Wave equ problem:	 Inference function - describing a solution to a specific data analytic system (a problem). For example, A linear (GLM) model represents a solution of a prediction inference problem, Y = Xβ, where the inference function quantifies the effects of all independent features (X) on the dependent outcome (Y), data: O = {X, Y}: ψ(O) = ψ(X,Y) ⇒ β̂ = β̂OLS = ⟨X X⟩⁻¹⟨X Y⟩ = (X^TX)⁻¹X^TY. 				
$ \begin{pmatrix} \frac{\partial^2}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2}{\partial t} \end{pmatrix} \psi(x,t) \\ = 0 $ Complex Solution: $ \psi(x,t) = A e^{i(kx - wt)} $ where $\left \frac{w}{k} \right = v$,	• A non-parametric, non-linear, alternative inference is SVM classification. If $\psi_x \in H$, is the lifting function $\psi \colon R^\eta \to R^d$ ($\psi \colon x \in R^\eta \to \tilde{x} = \psi_x \in H$), where $\eta \ll d$, the kernel $\psi_x(y) = \langle x y \rangle \colon 0 \times 0 \to R$ transformes non-linear to linear separation, the observed data $O_i = \{x_i, y_i\} \in R^\eta$ are lifted to $\psi_{O_i} \in H$. Then, the SVM prediction operator is the weighted sum of the kernel functions at ψ_{O_i} , where β^* is a solution to the SVM regularized optimization: $\langle \psi_O \mid \beta^* \rangle_H = \sum_{i=1}^n p_i^* \langle \psi_O \psi_{O_i} \rangle_H$				
represents a traveling wave	The linear coefficients, p_i^* , are the dual weights that are multiplied by the label corresponding to each training instance, $\{y_i\}$. Inference always depends on the (input) data; however, it does not have 1-1 and onto bijective correspondence with the data, since the inference function quantifies predictions in a probabilistic sense.				
GLM/S	VM: http://DSPA.predictive.space Dinov, Springer (2018)				

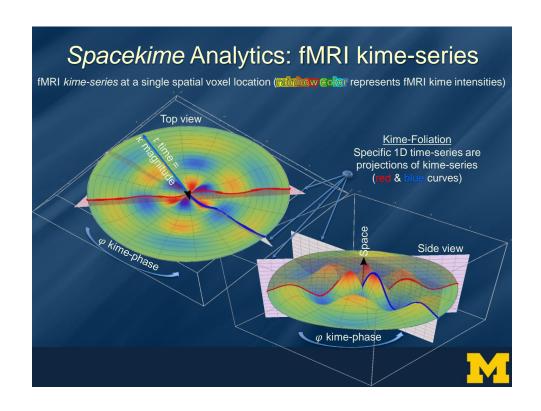


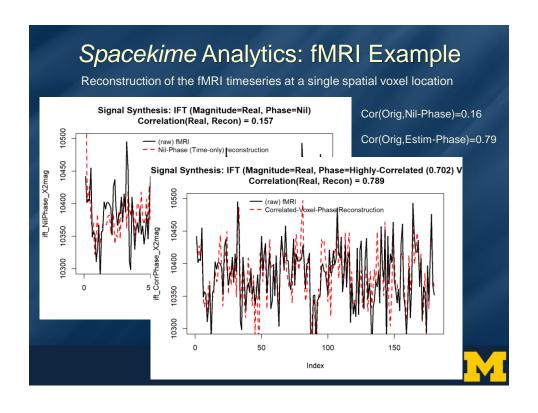


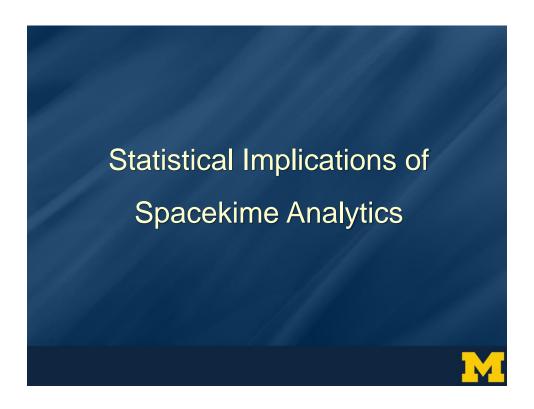












Uncertainty

- Quantum Mechanics: $||D_x u|| ||xu|| = \langle \frac{\hbar}{i} \partial_x u | ixu \rangle = \frac{\hbar}{2} ||u||^2 > 0$, i.e., non-commutation of the unbounded *operators* $D_x = \frac{\hbar}{i} \partial_x$ and x, (multiplication by x).
- Signal processing: Functions can't be time-limited **and** band-limited. Alternatively, a function and its Fourier transform cannot both have bounded domains $\sigma_t \times \sigma_\omega \geq 1/(4\pi)$, where σ_t, σ_ω are the time and frequency SDs.
- Intropic uncertainty: Entropy can be used just like the SD to quantify distribution structure. For instance, for angular, bimodal, or divergent-variance distributions, Entropy may be a better measure of dispersion than SD. For $FT(f)(\omega) = \hat{f}(\omega)$ and $IFT(\hat{f})(x) = \hat{f}(x)$, the Shannon information entropies:

$$H_x = \int \hat{f}(x) \log(\hat{f}(x)) dx$$
 and $H_\omega = \int \hat{f}(\omega) \log(\hat{f}(\omega)) d\omega$.

satisfy: $H_x + H_\omega \ge \log(e/2)$.

□ $L^2(\mathbb{R})$ <u>uncertainty</u>: it is impossible for $f \in L^2$ and \hat{f} to both decrease extremely rapidly. If both have rapidly decreasing tails: $|f(x)| \le C(1+|x|)^n e^{-a\pi x^2}$ and $|\hat{f}(\omega)| \le C(1+|\omega|)^n e^{-b\pi \omega^2}$, for some constant C, polynomial power n, and $a,b \in \mathbb{R}$, then f=0 (when ab>1); $f(x)=P_k(x)e^{-a\pi x^2}$ and $\hat{f}(\omega)=\widehat{P_k}(\omega)*e^{-\omega^2/4\pi a}$, where $\deg(P_k)\le n$ (when ab=1); or (when ab<1).



Heisenberg's Uncertainty in Spacekime?

- Heisenberg's uncertainty is resolved in 5D spacekime
- We can derive the classical 4D spacetime Heisenberg uncertainty as a reduction of Einstein-like 5D deterministic dynamics:
 - The math is terse it involves deriving the equations of motion by maximizing the distance (integral along the geodesic) between two points in 5D spacekime
 - The inner product du^{μ} $dx_{\mu}=\frac{dx^{\mu}dx_{\mu}}{L}=\frac{ds^2}{L}$. Since $\frac{ds}{L}\to 1$ near the leaf membrane, du^{μ} $dx_{\mu}=L=\frac{h}{mc}$. Replacing the change in velocity (du^{μ}) by the change in momentum (dp^{μ}) yields: $dp^{\mu} dx_{\mu} = h$.
 - This relation is similar to the quantum mechanics uncertainty principle in 4D Minkowski spacetime; however, it is obtained from 5D Einstein deterministic dynamics. In other words, in spacetime, Heisenberg's uncertainty principal manifests simply because of the one degree of freedom (kime-phase), i.e., lack of sufficient information about the second kime dimension.
 - In 5D spacekime, the conventional geodesic motion is perturbed by an extra force f^{μ} that can be split into two parts $f^{\mu}=f^{\mu}_{\perp}+f^{\mu}_{\parallel}$. The normal component f^{μ}_{\perp} is similar to other conventional forces and obeys the usual orthogonality condition $f_{\perp}^{\mu}u^{\mu}=0$. However, the parallel component f_{\parallel}^{μ} has no analog in 4D spacetime. In general, it has a non-trivial inner product with the 4-velocity u^{μ} , $f_{\parallel}^{\mu}u^{\mu}\neq 0$.
- In Minkowski 4D spacetime, the lack of kime-phase data naturally leaves one degree of freedom in the system causing Heisenberg's uncertainty. However, the latter can be explicated by information knowledge of the fifth component (kime-phase).

Wesson & Overduin, World Scientific (2018) | Dinov & Velev (2021)



Bayesian Inference Representation

- □ Suppose we have a single spacetime observation $X = \{x_{i_0}\} \sim p(x \mid \gamma)$ and $\gamma \sim$ $p(\gamma \mid \varphi = \text{phase})$ is a process parameter (or vector) that we are trying to estimate.
- \square Spacekime analytics aims to make appropriate inference about the process X.
- \square The <u>sampling distribution</u>, $p(x \mid \gamma)$, is the distribution of the observed data Xconditional on the parameter γ and the prior distribution, $p(\gamma \mid \varphi)$, of the parameter γ before the data X is observed, $\varphi = \text{phase aggregator}$.
- \square Assume that the hyperparameter (vector) φ , which represents the kime-phase estimates for the process, can be estimated by $\hat{\varphi} = \varphi'$.
- ☐ Such estimates may be obtained from an oracle, approximated using similar datasets, acquired as phases from samples of analogous processes, or derived via some phase-aggregation strategy.
- \square Let the <u>posterior distribution</u> of the parameter γ given the observed data $X = \{x_{i_o}\}$ be $p(\gamma|X, \varphi')$ and the process parameter distribution of the kime-phase hyperparameter vector φ be $\gamma \sim p(\gamma \mid \varphi)$.



Bayesian Inference Representation

☐ We can formulate spacekime inference as a Bayesian parameter estimation problem:

$$\underbrace{\frac{p(\gamma|X,\varphi')}{p(X|\varphi')}}_{\text{posterior distribution}} = \frac{\frac{p(\gamma,X,\varphi')}{p(X,\varphi')}}{\frac{p(X,\varphi')}{p(X|\varphi')}} = \frac{\frac{p(X|\gamma,\varphi')\times p(\gamma,\varphi')}{p(X|\varphi')\times p(\varphi')}}{\frac{p(X|\varphi')\times p(\varphi')}{p(X|\varphi')}} = \underbrace{\frac{p(X|\gamma,\varphi')\times p(\gamma,\varphi')}{p(X|\varphi')\times p(\varphi')}}_{\text{piserved evidence}} \propto \underbrace{\frac{p(X|\gamma,\varphi')\times p(\gamma,\varphi')}{p(X|\varphi')\times p(\gamma,\varphi')}\times \frac{p(\gamma,\varphi')}{p(\gamma,\varphi')}}_{\text{pisor}}.$$

- \Box In Bayesian terms, the posterior probability distribution of the unknown parameter γ is proportional to the product of the likelihood and the prior.
- \square In probability terms, the posterior = likelihood times prior, divided by the observed evidence, in this case, a single spacetime data point, x_{i_a} .



Bayesian Inference Representation

- \square Spacekime analytics based on a single spacetime observation x_{i_o} can be thought of as a type of Bayesian prior-predictive *or* posterior-predictive distribution estimation problem.
 - Prior predictive distribution of a new data point x_{j_o} , marginalized over the *prior* i.e., the sampling distribution $p(x_{j_o}|\gamma)$ weight-averaged by the pure *prior* distribution):

$$p(x_{j_o}|\varphi') = \int p(x_{j_o}|\gamma) \times \underbrace{p(\gamma|\varphi')}_{\text{prior distribution}} d\gamma.$$

□ Posterior predictive distribution of a new data point x_{i_0} , marginalized over the posterior; i.e., the sampling distribution $p(x_{i_0}|\gamma)$ weight-averaged by the posterior distribution:

$$p(x_{j_o}|x_{i_o},\varphi') = \int p(x_{j_o}|\gamma) \times \underbrace{p(\gamma|x_{i_o},\varphi')}_{\text{posterior distribution}} d\gamma.$$

- ☐ The difference between these two predictive distributions is that
 - \Box the posterior predictive distribution is updated by the observation $X = \{x_{i_o}\}$ and the hyperparameter, φ (phase aggregator),
 - □ whereas the prior predictive distribution only relies on the values of the hyperparameters that appear in the prior distribution.



Bayesian Inference Representation

- ☐ The posterior predictive distribution may be used to <u>sample</u> or <u>forecast</u> the distribution of a prospective, yet unobserved, data point x_{i_0} .
- □ The posterior predictive distribution spans the entire parameter state-space ($Domain(\gamma)$), just like the wavefunction represents the distribution of particle positions over the complete particle state-space.
- □ Using maximum likelihood or maximum *a posteriori* estimation, we can also estimate an individual parameter point-estimate, γ_o . In this frequentist approach, the point estimate may be plugged into the formula for the distribution of a data point, $p(x \mid \gamma_o)$, which enables drawing IID samples or individual outcome values.



Bayesian Inference Simulation

- □ Simulation example using 2 random samples drawn from mixture distributions each of $n_A = n_B = 10$ K observations:
 - \square $\{X_{A,i}\}_{i=1}^{n_A}$, where $X_{A,i} = 0.3U_i + 0.7V_i$, $U_i \sim N(0,1)$ and $V_i \sim N(5,3)$, and
 - \square $\{X_{B,i}\}_{i=1}^{n_B}$, where $X_{B,i}=0.4P_i+0.6Q_i, P_i\sim N(20,20)$ and $Q_i\sim N(100,30)$.
- □ The intensities of cohorts A and B are independent and follow different mixture distributions. We'll split the first cohort (A) into training (C) and testing (D) subgroups, and then:
 - ☐ Transform all four cohorts into Fourier k-space,
 - \Box Iteratively randomly sample single observations from (training) cohort C,
 - \square Reconstruct the data into spacetime using a single kime-magnitude value and alternative kime-phase estimates derived from cohorts B, C, and D, and
 - ☐ Compute the classical spacetime-derived population characteristics of cohort *A* and compare them to their spacekime counterparts obtained using a single *C* kime-magnitude paired with *B*, *C*, or *D* kime-phases.





values of their counterparts computed using the spacekime reconstructed signals based on kime-phases of cohorts B, C, and D. The <u>estimates for the latter three cohorts correspond to reconstructions using a single spacetime observation (i.e., single kime-magnitude) and alternative kime-phases (in this case, kime-phases derived from cohorts B, C, and D).</u>

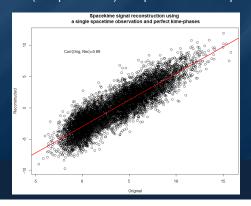
		Spacetime Spacekime Reconstructions (single kime-magnitu				
	C	(A)	(B)	(<i>C</i>)	(D)	
	Summaries	Original	Phase=Diff. Process	Phase=True	Phase=Independent	
	Min	-2.38798	-3.798440	-2.98116	-2.69808	
	1 st Quartile	-0.89359	-0.636799	-0.76765	-0.76453	
	Median	0.03311	0.009279	-0.05982	-0.08329	
	Mean	0.00000	0.000000	0.00000	0.00000	
	3 rd Quartile	0.75772	0.645119	0.72795	0.69889	
	Max	3.61346	3.986702	3.64800	3.22987	
1	Skewness	0.348269	0.001021943	0.2372526	0.31398	
0.10-	Kurtosis	-0.68176	0.2149918	-0.4452207	-0.3270084	
density	cohort			100		



Bayesian Inference Simulation

The correlation between the original data (*A*) and its <u>reconstruction using a single kime magnitude</u> and the correct kime-phases (*C*) is $\rho(A, C) = 0.89$.

This strong correlation suggests that a substantial part of the A process <u>energy</u> <u>can be recovered using only a single observation</u>. In this case, to reconstruct the signal back into spacetime and compute the corresponding correlation, we used a single kime-magnitude (sample-size=1) and process \mathcal{C} kime-phases.





Bayesian Inference Simulation

Let's demonstrate the Bayesian inference corresponding to this spacekime data analytic problem using a simulated bimodal experiment:

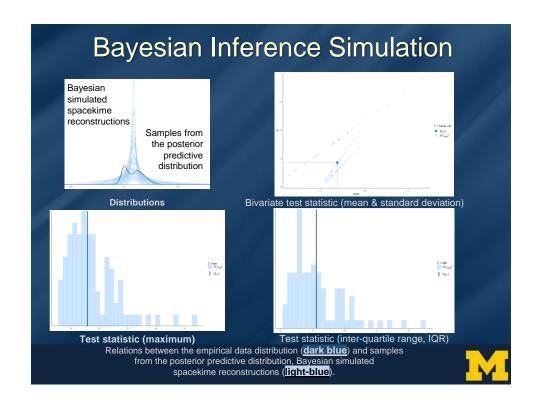
 $\overline{X_A} = 0.3U + 0.7V$, where $U \sim N(0,1)$ and $V \sim N(5,3)$

Specifically, we will illustrate the Bayesian inference using repeated single spacetime observations from cohort $A, X = \{x_{i_o}\}$, and varying kime-phase priors $(\theta = \text{phase aggregator})$ obtained from cohorts B, C, or D, using different posterior predictive distributions

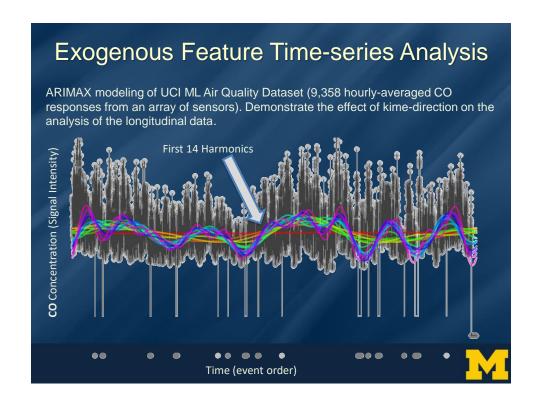
Relations between the empirical data distribution (<u>dark blue</u>) and samples from the posterior predictive distribution, representing Bayesian simulated spacekime reconstructions (<u>light-blue</u>). The derived Bayesian estimates do not perfectly match the empirical distribution of the simulated data, yet there is clearly information encoding that is captured by the spacekime data reconstructions

This <u>signal compression</u> can be exploited by subsequent model-based or model-free data analytic strategies for retrospective prediction, prospective forecasting, ML classification, derived clustering, and other spacekime inference methods









Exo	genous I	Feature	e Time-s	series An	alysis
	Phase	Nil	Average	True=original	
0	Model Estimate	ARIMA(2,0,1)	ARIMA(2,0,3)	ARIMA(1,1,4)	9 0 0
3 different alternative	AIC	13179	14183	10581	
er(ar1	1.11406562	0.329482302	0.2765312	ARIM order (# of time ladifferencing (# of order of MA part
diff ter	ar2	-0.14565048	0.238363531		er (er e
	ma1	-0.78919188	0.267291585	-0.88913497	(# c enci
on	ma2		-0.006079386	0.12679494	of t
	ma3		0.15726556	0.03043726	ARIN time I g (# of A part
	ma4			-0.17655728	
	intercept	503.3455144	742.800113		ARIMAX time lags) g (# of pas A part
ls (ons ma	xreg1	-0.40283891	0.58379483	0.08035744	st v of 6
models tructions n estima	xreg2	0.13656613	0.280936931	6.14947902	(p,d,q) of the AR it values si
e ii i	xreg3	-0.51457636	-0.649722755	0.09859223	e / e /
X r nst	xreg4	1.09611981	1.239910298	0.01634736	su R
AA Sct	xreg5	1.21946209	-0.026110332	-0.04816591	part
CO ARIMAX models derived signal reconstructions based kime-direction estimates	xreg6	1.30628469	1.081777956	-0.01104142	rt
	xreg7	1.20868397	0.254018471	0.1832854	AX (p,d,q) ags) of the AR part past values subtractions
	xreg8	1.14905809	0.306524131	0.17648482	15)
	xreg9	-0.48233756	-0.405204908	6.53739782	
	xreg10	0.03145281	0.351063312	1.79388326	-
	xreg11	-0.46395772	-0.457689796	-12.06965578	

